# SEARCHES FOR MONOPOLES, SUPERSYMMETRY, TECHNICOLOR, COMPOSITENESS, EXTRA DIMENSIONS, etc.

## Magnetic Monopole Searches

Isolated supermassive monopole candidate events have not been confirmed. The most sensitive experiments obtain negative results.

Best cosmic-ray supermassive monopole flux limit:

$$< 1.4 \times 10^{-16} \text{ cm}^{-2} \text{sr}^{-1} \text{s}^{-1}$$
 for  $1.1 \times 10^{-4} < \beta < 1$ 

# Supersymmetric Particle Searches

Presently all supersymmetric mass bounds are model dependent.

This table contains a selection of bounds indicating the range of possibilities. For a more extensive set of cases consult the detailed listings.

The limits are based on the Minimal Supersymmetric Standard Model (MSSM) with additional assumptions as follows:

1)  $\tilde{\chi}_1^0$  is lightest supersymmetric particle; 2) *R*-parity is conserved;

See the Particle Listings for a Note giving details of supersymmetry.

$$\begin{array}{l} \widetilde{\chi}_{i}^{0} - \text{neutralinos (mixtures of } \widetilde{\gamma}, \ \widetilde{Z}^{0}, \ \text{and } \widetilde{H}_{i}^{0}) \\ \text{Mass } m_{\widetilde{\chi}_{1}^{0}} > 0 \ \text{GeV, CL} = 95\% \\ \text{[general MSSM, non-universal gaugino masses]} \\ \text{Mass } m_{\widetilde{\chi}_{1}^{0}} > 46 \ \text{GeV, CL} = 95\% \\ \text{[all } \tan\beta, \ \text{all } m_{0}, \ \text{all } m_{\widetilde{\chi}_{2}^{0}} - m_{\widetilde{\chi}_{1}^{0}}] \\ \text{Mass } m_{\widetilde{\chi}_{2}^{0}} > 62.4 \ \text{GeV, CL} = 95\% \\ \text{[$1$<$} \tan\beta < 40, \ \text{all } m_{0}, \ \text{all } m_{\widetilde{\chi}_{2}^{0}} - m_{\widetilde{\chi}_{1}^{0}}] \\ \text{Mass } m_{\widetilde{\chi}_{2}^{0}} > 345 \ \text{GeV, CL} = 95\% \\ \text{[$\widetilde{\chi}_{1}^{\pm} \ \widetilde{\chi}_{2}^{0} \rightarrow \ W \ \widetilde{\chi}_{1}^{0} \ Z \ \widetilde{\chi}_{1}^{0}, \ \text{simplified model, } m_{\widetilde{\chi}_{1}^{\pm}} = m_{\widetilde{\chi}_{2}^{0}}, \ m_{\widetilde{\chi}_{1}^{0}} = 0 \ \text{GeV}] \\ \end{array}$$

Citation: C. Patrignani et al. (Particle Data Group), Chin. Phys. C, 40, 100001 (2016)

Mass 
$$m_{\widetilde{\chi}^0_3} > 99.9$$
 GeV,  $\mathsf{CL} = 95\%$ 

$$m_{\widetilde{\chi}_3^0} > 35.3 \text{ GeV}, \text{ CL} = 33.6$$
 [1 $<$ tan $eta <$ 40, all  $m_0$ , all  $m_{\widetilde{\chi}_2^0} - m_{\widetilde{\chi}_1^0}$ ]

Mass 
$$m_{\widetilde{\chi}_4^0} > 116$$
 GeV, CL = 95%

[1
$$<$$
tan $eta$   $<$ 40, all  $m_0$ , all  $m_{\widetilde{\chi}^0_2} - m_{\widetilde{\chi}^0_1}$ ]

$$\widetilde{\chi}_i^\pm$$
 — charginos (mixtures of  $\widetilde{W}^\pm$  and  $\widetilde{H}_i^\pm$ )

Mass  $m_{\widetilde{\chi}_1^\pm} > 94$  GeV, CL  $= 95\%$ 

$$[ aneta<40,\ m_{\widetilde{\chi}_1^\pm}-m_{\widetilde{\chi}_1^0}>3\ ext{GeV},\ ext{all}\ m_0]$$

Mass 
$$m_{\widetilde{\chi}_1^{\pm}} > 345$$
 GeV, CL = 95%

[simplified model, 
$$m_{\widetilde{\chi}_1^\pm}=m_{\widetilde{\chi}_2^0}$$
,  $m_{\widetilde{\chi}_1^0}=0$  GeV]

$$\widetilde{\nu}$$
 — sneutrino

Mass 
$$m>94$$
 GeV, CL  $=95\%$  [CMSSM,  $1\leq \tan\beta \leq 40$ ,  $m_{\widetilde{e}_R}-m_{\widetilde{\chi}_1^0}>10$  GeV]

[CMSSM, 
$$1 \leq \tan \beta \leq 40$$
,  $m_{\widetilde{e}_R} - m_{\widetilde{\chi}_1^0} > 10$  GeV]

$$\widetilde{e}$$
 — scalar electron (selectron)

Mass 
$$m(\widetilde{e}_L)>107$$
 GeV, CL  $=95\%$  [all  $m_{\widetilde{e}_R}-m_{\widetilde{\chi}_1^0}$ ] Mass  $m(\widetilde{e}_R)>97.5$  GeV, CL  $=95\%$  [ $\Delta m>11$  GeV,  $|\mu|>100$  GeV,  $\tan\beta=1.5$ ]

$$\widetilde{\mu}$$
 — scalar muon (smuon)  
Mass  $m>~94$  GeV, CL  $=95\%$ 

[CMSSM, 
$$1 \leq aneta \leq 40$$
,  $m_{\widetilde{\mu}_R} - m_{\widetilde{\chi}^0_1} > 10$  GeV]

$$\widetilde{ au}$$
 — scalar tau (stau)  
Mass  $m>81.9$  GeV, CL  $=95\%$ 

$$[m_{\widetilde au_R}-m_{\widetilde\chi^0_1}$$
 >15 GeV, all  $heta_ au$ , B $(\widetilde au o$   $au\,\widetilde\chi^0_1)=$  100%]

$$\chi_1^0$$
 > 25 331, 21  $\delta_{\gamma}$ , 2(...

$$\widetilde{q}$$
 – squarks of the first two quark generations

The first of these limits is within CMSSM with cascade decays, evaluated assuming a fixed value of the parameters 
$$\mu$$
 and  $\tan \beta$ . The first two limits assume two-generations of mass degenerate squarks ( $\tilde{q}_{\mu}$  and  $\tilde{q}_{\mu}$ ) and gaus

assume two-generations of mass degenerate squarks  $(\widetilde{q}_L \text{ and } \widetilde{q}_R)$  and gaugino mass parameters that are constrained by the unification condition at the grand unification scale. The third limit assumes a simplified model with a

100% branching ratio for the prompt decay  $\widetilde{q} \rightarrow q \widetilde{\chi}_1^0$ . Mass m > 1450 GeV, CL = 95%

[CMSSM, 
$$tan\beta = 30$$
,  $A_0 = -2max(m_0, m_{1/2})$ ,  $\mu > 0$ ]

Mass 
$$m>850$$
 GeV, CL  $=95\%$   
[jets  $+\cancel{E}_T$ ,  $\widetilde{q}\to q\widetilde{\chi}^0_1$  simplified model,  $m_{\widetilde{\chi}^0_1}=0$  GeV]

Mass 
$$m>520$$
 GeV, CL  $=95\%$   $[\widetilde{q}\to q\widetilde{\chi}^0_1, \text{ simplified model, single light squark, } m_{\widetilde{\chi}^0_1}=0]$ 

$$\widetilde{b}$$
 — scalar bottom (sbottom)

Mass  $m > 650$  GeV, CL = 95%  $[\widetilde{b} \rightarrow b\widetilde{\chi}_1^0, m_{\widetilde{\chi}_1^0} = 0]$ 

Mass  $m > 600$  GeV, CL = 95%  $[\widetilde{b} \rightarrow b\widetilde{\chi}_1^0, m_{\widetilde{\chi}_1^0} < 250$  GeV]

 $\widetilde{t}$  — scalar top (stop)

Mass  $m > 730$  GeV, CL = 95%

$$\begin{split} & [\widetilde{t} \to t \, \widetilde{\chi}_1^0, \, m_{\widetilde{\chi}_1^0} = 100 \, \text{GeV}, \, m_{\widetilde{t}} > m_t + m_{\widetilde{\chi}_1^0}] \\ & \text{Mass } m > 500 \, \text{GeV}, \, \text{CL} = 95\% \\ & [\ell^\pm + \text{jets} + \not\!\!E_T, \, \widetilde{t}_1 \to \, b \, \widetilde{\chi}_1^\pm, \, m_{\widetilde{\chi}_1^\pm} = 2 \, m_{\widetilde{\chi}_1^0}, \, 100 \, \, \text{GeV} < m_{\widetilde{\chi}_1^0} \, < 150 \end{split}$$

GeV] Mass 
$$m > 240$$
 GeV,  $CL = 95\%$ 

Mass 
$$m > 240$$
 GeV,  $CL = 95\%$   

$$[\widetilde{t}_1 \rightarrow c \widetilde{\chi}_1^0, m_{\widetilde{t}_1} - m_{\widetilde{\chi}_1^0} < 85 \text{ GeV}]$$

$$\widetilde{g}$$
 — gluino

the prompt 3 body decay, independent of the squark mass. The second of these limits is within the CMSSM (for  $m_{\widetilde{g}} \gtrsim 5$  GeV), and includes the effects of cascade decays, evaluated assuming a fixed value of the parameters  $\mu$  and  $tan\beta$ . The limit assumes GUT relations between gaugino masses and the gauge couplings. The third limit is based on a combination of searches.

The first limit assumes a simplified model with a 100% branching ratio for

Mass 
$$m>1225$$
 GeV, CL  $=95\%$   $[\widetilde{g}\to q\overline{q}\widetilde{\chi}_1^0,\ m_{\widetilde{\chi}_1^0}=0]$  Mass  $m>1150$  GeV, CL  $=95\%$ 

[CMSSM, 
$$\tan\beta$$
=30,  $A_0$ = $-2\max(m_0,m_{1/2})$ ,  $\mu>0$ ] Mass  $m>1150$  GeV, CL = 95%

[general RPC 
$$\tilde{g}$$
 decays,  $m_{\approx 0}$  < 100 GeV

[general RPC  $\widetilde{g}$  decays,  $m_{\widetilde{\chi}_1^0}$  < 100 GeV]

## **Technicolor**

The limits for technicolor (and top-color) particles are quite varied depending on assumptions. See the Technicolor section of the full Review (the data listings).

# Quark and Lepton Compositeness, Searches for

# Scale Limits $\Lambda$ for Contact Interactions (the lowest dimensional interactions with four fermions)

If the Lagrangian has the form

$$\pm \frac{g^2}{2\Lambda^2} \overline{\psi}_{\mathsf{L}} \gamma_{\mu} \psi_{\mathsf{L}} \overline{\psi}_{\mathsf{L}} \gamma^{\mu} \psi_{\mathsf{L}}$$

 $\Lambda_{II}^{+}(eeee) > 8.3 \text{ TeV}, CL = 95\%$ 

(with  $g^2/4\pi$  set equal to 1), then we define  $\Lambda \equiv \Lambda_{LL}^{\pm}$ . For the full definitions and for other forms, see the Note in the Listings on Searches for Quark and Lepton Compositeness in the full *Review* and the original literature.

$$\Lambda_{LL}^{-}(eeee)$$
 > 10.3 TeV, CL = 95%  $\Lambda_{LL}^{+}(ee\mu\mu)$  > 8.5 TeV, CL = 95%  $\Lambda_{LL}^{-}(ee\mu\mu)$  > 9.5 TeV, CL = 95%  $\Lambda_{LL}^{-}(ee\tau\tau)$  > 7.9 TeV, CL = 95%  $\Lambda_{LL}^{-}(ee\tau\tau)$  > 7.2 TeV, CL = 95%  $\Lambda_{LL}^{-}(ee\tau\tau)$  > 9.1 TeV, CL = 95%  $\Lambda_{LL}^{-}(\ell\ell\ell\ell)$  > 9.1 TeV, CL = 95%  $\Lambda_{LL}^{-}(\ell\ell\ell\ell)$  > 10.3 TeV, CL = 95%  $\Lambda_{LL}^{-}(eeuu)$  > 23.3 TeV, CL = 95%  $\Lambda_{LL}^{-}(eeuu)$  > 12.5 TeV, CL = 95%  $\Lambda_{LL}^{-}(eedd)$  > 11.1 TeV, CL = 95%  $\Lambda_{LL}^{-}(eedd)$  > 26.4 TeV, CL = 95%  $\Lambda_{LL}^{-}(eecc)$  > 9.4 TeV, CL = 95%  $\Lambda_{LL}^{-}(eebb)$  > 9.4 TeV, CL = 95%  $\Lambda_{LL}^{-}(eebb)$  > 10.2 TeV, CL = 95%  $\Lambda_{LL}^{-}(eebb)$  > 10.2 TeV, CL = 95%  $\Lambda_{LL}^{-}(\mu\mu qq)$  > 16.7 TeV, CL = 95%  $\Lambda_{LL}^{-}(\mu\mu qq)$  > 16.7 TeV, CL = 95%  $\Lambda_{LL}^{-}(\mu\mu qq)$  > 2.81 TeV, CL = 95%  $\Lambda_{LL}^{-}(qqqq)$  > 2.81 TeV, CL = 95%  $\Lambda_{LL}^{-}(qqqq)$  > 12.0 TeV, CL = 95%  $\Lambda_{LL}^{-}(qqqq)$  > 12.0 TeV, CL = 95%  $\Lambda_{LL}^{-}(qqqq)$  > 12.0 TeV, CL = 95%  $\Lambda_{LL}^{-}(\mu\nu qq)$  > 5.0 TeV, CL = 95%  $\Lambda_{LL}^{-}(\nu\nu qq)$  > 5.4 TeV, CL = 95%

#### **Excited Leptons**

The limits from  $\ell^{*+}\ell^{*-}$  do not depend on  $\lambda$  (where  $\lambda$  is the  $\ell\ell^{*}$  transition coupling). The  $\lambda$ -dependent limits assume chiral coupling.

$$e^{*\pm}$$
 — excited electron

Mass 
$$m > 103.2 \text{ GeV}$$
,  $CL = 95\%$  (from  $e^* e^*$ )

Mass 
$$m > 3.000 \times 10^3$$
 GeV, CL = 95% (from  $ee^*$ )  
Mass  $m > 356$  GeV, CL = 95% (if  $\lambda_{\gamma} = 1$ )

$$\mu^{*\pm}$$
 — excited muon

Mass 
$$m>103.2$$
 GeV, CL = 95% (from  $\mu^*\mu^*$ )

Mass 
$$m>~3.000 imes10^3$$
 GeV, CL  $=95\%$  (from  $\mu\mu^*$ )

$$au^{*\pm}$$
 — excited tau

Mass 
$$m > 103.2$$
 GeV, CL = 95% (from  $\tau^* \tau^*$ )  
Mass  $m > 2.500 \times 10^3$  GeV, CL = 95% (from  $\tau \tau^*$ )

$$\nu^*$$
 — excited neutrino

Mass 
$$m>1.600\times 10^3$$
 GeV, CL  $=95\%$  (from  $\nu^*\nu^*$ )

Mass 
$$m > 1.000 \times 10^{\circ}$$
 GeV,  $CL = 95\%$  (from  $\nu^* X$ )

Mass 
$$m > 338 \text{ GeV}$$
,  $CL = 95\%$  (from  $q^* q^*$ )

Mass 
$$m > 4.060 \times 10^3$$
 GeV,  $CL = 95\%$  (from  $q^* X$ )

#### Color Sextet and Octet Particles

Color Sextet Quarks 
$$(q_6)$$

Mass 
$$m>84$$
 GeV,  $CL=95\%$  (Stable  $q_6$ )

Color Octet Charged Leptons 
$$(\ell_8)$$

Mass 
$$m > 86$$
 GeV,  $CL = 95\%$  (Stable  $\ell_8$ )

Color Octet Neutrinos (
$$\nu_8$$
)

Mass 
$$m>~110$$
 GeV,  $\mathsf{CL}=90\%~~(
u_8 
ightarrow~
u_g)$ 

## **Extra Dimensions**

Please refer to the Extra Dimensions section of the full *Review* for a discussion of the model-dependence of these bounds, and further constraints.

# Constraints on the radius of the extra dimensions, for the case of two-flat dimensions of equal radii

$$R < 30 \ \mu \text{m}$$
, CL = 95% (direct tests of Newton's law)

$$R < 15 \ \mu \text{m}, \text{ CL} = 95\% \quad (pp \rightarrow jG)$$

$$R < 0.16-916$$
 nm (astrophysics; limits depend on technique and assumptions)

#### Constraints on the fundamental gravity scale

$$M_{TT}>$$
 6.3 TeV, CL  $=95\%$  (pp  $\to$  dijet, angular distribution)  $M_c>$  4.16 TeV, CL  $=95\%$  (pp  $\to$   $\ell \overline{\ell})$ 

### Constraints on the Kaluza-Klein graviton in warped extra dimensions

$$M_G~>~2.73$$
 TeV, CL  $=95\%~~(pp
ightarrow~e^+e^-,~\mu^+\mu^-)$ 

#### Constraints on the Kaluza-Klein gluon in warped extra dimensions

$$\mathit{M}_{\mathit{g}_{KK}}~>~2.5$$
 TeV, CL  $=95\%~~(\mathit{g}_{KK} 
ightarrow~t\,\overline{t})$